

# The Influence of Ion-Implanted Profiles on the Performance of GaAs MESFET's and MMIC Amplifiers

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**Abstract** — The RF small-signal performance of GaAs MESFET's and MMIC amplifiers as a function of various ion-implanted profiles is theoretically and experimentally investigated. Implantation energy, dose, and recess depth influence are theoretically analyzed with the help of a specially developed device simulator. The performance of MMIC amplifiers processed with various energies, doses, recess depths, and bias conditions is discussed and compared to experimental characteristics. Some criteria are finally proposed for the choice of implantation conditions and process in order to optimize the characteristics of ion-implanted FET's and to realize process-tolerant MMIC amplifiers.

## I. INTRODUCTION

THE ELECTRICAL performance of GaAs MESFET's is greatly influenced by the doping profile of their active layer [1]–[3]. In addition, the rapid development of active layer fabrication techniques (VPE, MBE, ion implantation) allows the realization of almost any doping profile. Therefore, for discrete devices or monolithic integrated circuits, the doping has to be specially tailored to provide the best microwave performance.

In this paper, we analyze the influence of ion implantation conditions (dose, energy) and recess depth on the main elements of the MESFET equivalent circuit and on the microwave small-signal performance of ion-implanted MMIC amplifiers.

Section II describes the simulator used for our investigations and discusses some theoretical results obtained with it.

In Section III, the MMIC amplifier design is described together with the influence that dose, energy, and recess have on its performance. Finally, Section IV compares data to theoretical predictions for MESFET's and the MMIC amplifier.

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## II. THE SIMULATOR "SIMTEC"

### A. Model Description

Several models for nonuniformly doped MESFET's have been reported [4]–[7]. The simulator described in this paper was developed in order to satisfy the following requirements: (i) to investigate any doping profile without the limitations imposed by equivalent profile transformations; (ii) to simulate submicrometer gate devices with short computation time (the influence of gate length on velocity is considered without performing long two-dimensional computations); and (iii) to compute all the elements of the equivalent circuit, including, for instance, the input gate resistance  $R_i$ , which is often neglected but strongly affects  $f_{\max}$ .

The simulations presented below are based on a "hybrid" one-dimensional model involving an analytical formalism and numerical treatments of the classical equations of the "two-region" model [8]. In this way, it is possible to predict the performance of a GaAs MESFET with any doping profile described by a mathematical expression [7]. This model is, therefore, particularly suitable for ion-implanted MESFET's, where the doping profile shape depends on various parameters such as energy and dose of implantation.

A first description of this new simulator, called SIMTEC, has already been presented by the authors elsewhere [9]. From the MESFET geometrical dimensions and physical parameters given in Fig. 1 and Table I, the simulator SIMTEC yields for any doping profile and bias conditions  $V_{DS}$ ,  $V_{GS}$  (Or  $I_{DS}$ ):

- all the elements of the MESFET equivalent circuit except the series inductances (see Fig. 2 by way of example);
- the values of the current cutoff frequency  $f_T$ , the maximum oscillation frequency  $f_{\max}$ , and the minimum noise figure  $F_{\min}$ .

The model is suitable for submicrometer gate devices, the influence of velocity overshoot being taken into account by the following approximate relationship, proposed in

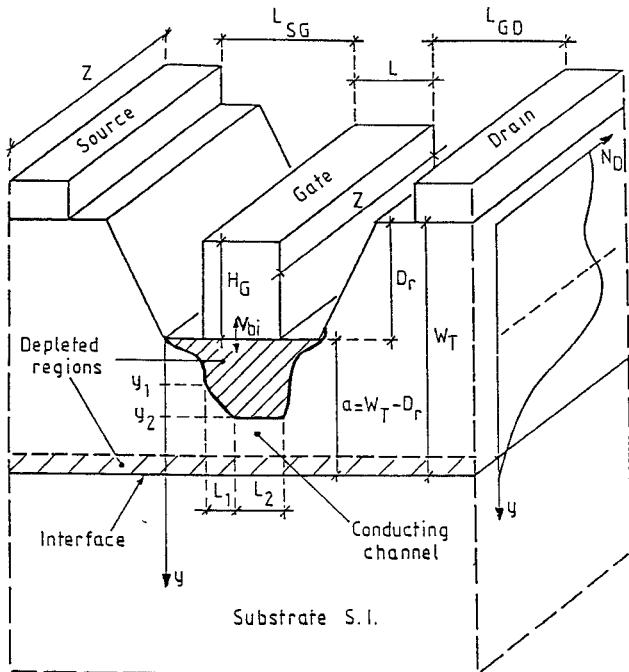


Fig. 1. Idealized cross section of MESFET devices used in this work.

TABLE I  
GEOMETRICAL AND ELECTRICAL PARAMETERS USED IN  
THE MODEL (SIMTEC)

$L$	gate length
$Z$	total gate width
$L_{SG}$	source-gate spacing
$L_{GD}$	gate-drain spacing
$D_r$	recess depth
$\rho_c$	ohmic contact specific resistance
$V_{bi}$	Schottky barrier built-in potential
$H_G$	gate metallization height
$N_k$	number of gate fingers of width $z$ ( $Z = N_k \cdot z$ )
$\rho_G$	gate metal resistivity
$R_{TH}$	thermal resistance

The values of  $L_{SG}$ ,  $L_{GD}$ ,  $D_r$ ,  $\rho_c$ ,  $H_G$ ,  $\rho_G$ , and  $N_k$  make it possible to determine the access resistances  $R_S$ ,  $R_D$ , and  $R_G$ .

[10]:

$$v_{\text{sat}} = 60L^{-0.56} \quad (1)$$

where  $L$  is the gate length (m) and  $v_{\text{sat}}$  the equivalent saturated velocity (m/s). This relationship gives a realistic effective saturation velocity  $v_{\text{sat}}$ , which is larger than the steady-state velocity. Thus  $v_{\text{sat}}$  is  $1.4 \times 10^5$  m/s if  $L = 1 \mu\text{m}$  and reaches the effective value of  $2.7 \times 10^5$  m/s if  $L = 0.3 \mu\text{m}$ .

For reduced computation times, the mobility is assumed to have a constant average value  $\mu_0$  across the active layer. SIMTEC can, however, accept precise functions  $\mu_0(y)$  for the mobility profile. The influence of ionized impurities on the average mobility is approximated by

$$\mu_0(\text{cm}^2\text{V}^{-1}\text{s}^{-1}) = \frac{8000}{1 + \left[ \frac{Neq}{10^{17}} \right]^{1/2}} \quad (2a)$$

where  $Neq$  ( $\text{cm}^{-3}$ ) is the average impurity concentration,

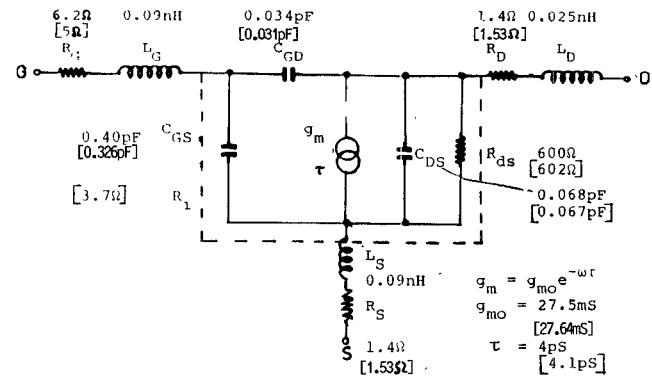


Fig. 2. FET equivalent circuit de-embedded from S-parameter measurements and predicted by theory (in brackets). Implantation conditions:  $E_1 = 60 \text{ keV}$ ,  $D_1 = 20 \times 10^{12} \text{ cm}^{-2}$ , and  $E_2 = 250 \text{ keV}$ ,  $D_2 = 6 \times 10^{12} \text{ cm}^{-2}$ . Bias point:  $V_{DS} = 3 \text{ V}$ ,  $I_{DS} = 30 \text{ mA}$ . A recess depth of  $0.14 \mu\text{m}$  provides  $V_T = -3.6 \text{ V}$ .

given by

$$Neq = \frac{\left[ \int_0^\infty N_D(y) dy \right]^2}{2 \int_0^\infty N_D(y) y dy} \quad (2b)$$

$N_D(y)$  being the doping density. Therefore the mobility degradation for the highest doping densities [11] is systematically considered in our simulations.

Two-piece velocity-electric field characteristics are employed using the following relations:

$$v(E) = \frac{\mu_0 E}{1 + E/E_0} \quad \text{if } E \leq E_s \quad (2c)$$

and

$$v(E) = v_{\text{sat}} \quad \text{if } E > E_s \quad (2d)$$

where  $E_0$  is a parameter with electric field dimensions equal to  $2.5 \times 10^6 \text{ V m}^{-1}$ .  $E_s$  is the saturation electric field ( $E_s = v_{\text{sat}}/(\mu_0 - v_{\text{sat}}/E_0)$ ).

The simulator has two additional features. The first is that the Debye transition zone between the depletion layer and the conducting channel can be taken into account in the computations. The model can indeed distinguish optionally between free electron and shallow level impurity concentrations. By solving the semiconductor equations at all points of the channel, one can remove the limitation of abrupt transition between the depleted and neutral part of the channel. This results in a precise calculation of the free carrier distribution  $n(y)$ , which over several Debye lengths  $\lambda_D = [\epsilon KT/q^2 N_D(y)]^{1/2}$  is equal to neither zero nor  $N_D(y)$ .

The second additional feature is that the presence in the semi-insulating substrate of deep levels ( $EL_2$ , for example) forming an N-I junction with the N-doped layer of the MESFET [12] can also be accounted for.

These optional capabilities were not retained in this work in order to simplify the analysis and minimize computation time. At given bias conditions, the execution time could therefore be less than 0.1 s for dc and small-signal simulation on an IBM 3080 computer. This is many orders

of magnitude less than a numerical simulation [13]. Therefore, SIMTEC is very suitable for the computer-aided design of low-complexity circuits.

### B. Method of Element Calculation

The basic device structure used in the simulations is shown in Fig. 1. The technology parameters of Table II were maintained constant throughout this work and correspond to the fabricated devices. For every implantation condition (energy, dose) or recess depth, all the elements of the model (Fig. 2) are calculated, with the exception of  $L_S$ ,  $L_D$ , and  $L_G$ . The computation is based on the solution of the following system of equations.

1) *Linear (Ohmic) Operation Equations*: These are derived from Poisson's equations by considering the current  $I_{DS}$  across the channel and the drain-source  $V_{DS}$  and gate-source  $V_{GS}$  voltages:

$$V_{bi} + R_S I_{DS} - V_{GS} = \frac{q}{\epsilon} \int_0^{y_1} y N_D(y) dy \quad (3)$$

$$V_{DS} - (R_S + R_D) I_{DS} = \frac{q}{\epsilon} \int_{y_1}^{y_2} y N_D(y) dy \quad (4)$$

$$\begin{aligned} I_{DS} \left( L + \frac{q}{\epsilon E_0} \int_{y_1}^{y_2} y N_D(y) dy \right) \\ = \frac{q^2 Z}{\epsilon} \int_{y_1}^{y_2} \left( \int_y^a \mu_0(y') N_D(y') dy' \right) y N_D(y) dy. \end{aligned} \quad (5)$$

2) *Saturation Operation Equations*: The two-region model [8] is employed. Equations (3) and (5) can be used again in linear region I by replacing  $L$  with  $L_1$ . In region II carrier propagation is at the saturated velocity  $v_{sat}$  and the following equation holds:

$$I_{DS} = q v_{sat} Z \int_{y_2}^a N_D(y) dy. \quad (6)$$

Using the method of Grebene and Ghandi [14] for the solution of the two-dimensional Poisson's equation and defining an equivalent channel thickness  $a_s$  as

$$a_s = \frac{I_{DS}}{q v_{sat} Z N_D(y_2)} + y_2 \quad (7)$$

we obtain

$$\begin{aligned} V_{DS} - (R_S + R_D) I_{DS} = \frac{q}{\epsilon} \int_{y_1}^{y_2} y N_D(y) dy \\ + \frac{2 a_s E_s}{\pi} \sinh \left( \frac{\pi(L - L_1)}{2 a_s} \right). \end{aligned} \quad (8)$$

Depending on the operating conditions  $V_{DS}$ ,  $V_{GS}$ , the static parameters of the MESFET can be obtained by numerical solution of either (3)–(5) (evaluation of  $y_1$ ,  $y_2$ ,  $I_{DS}$  in linear mode) or (3), (5)–(8) (evaluation of  $a_s$ ,  $y_1$ ,  $y_2$ ,  $L_1$ ,  $I_{DS}$  in saturation mode). The equivalent circuit parameters  $g_{m0}$ ,  $R_{DS}$  are obtained from the ratio of the incremental change of drain current  $\Delta I_{DS}$  to the incremental change of

TABLE II  
CONSTANT PARAMETERS USED IN THE SIMULATION  
IRRESPECTIVE OF DOPING PROFILE

$L$	1 $\mu\text{m}$
$Z$	300 $\mu\text{m}$
$L_{SG}$	1.75 $\mu\text{m}$
$L_{GD}$	1.75 $\mu\text{m}$
$V_{bi}$	0.75 V

The significance of  $L$ ,  $Z$ ,  $L_{SG}$ , and  $L_{GD}$  is given in Fig. 1.  $V_{bi}$  is the built-in gate voltage.

intrinsic drain  $\Delta V_{DS}$  and gate voltage  $\Delta V_{GS}$ .  $C_{GS}$  is evaluated using the transmission line approach described in the Appendix. These calculations can be performed with the help of the static equations given above. For the exact calculation of  $R_{DS}$ , we included the effects of carrier injection into the semi-insulating substrate by considering a parallel connected parasitic resistance  $R_{DS_p}$ . This resistance is found to be proportional to the length  $L - L_1$  of the saturated region [9].

The access inductances  $L_S$ ,  $L_D$ ,  $L_G$  are not calculated with the program and were extracted from experimental  $S$ -parameter data.

The source and drain series resistances  $R_S$  and  $R_D$  consist of a contact resistance  $R_{SC}$  in series with the lateral active layer resistance  $R_{SL}$ .  $R_{SC}$  is a function of the specific contact resistivity  $\rho_c$ , which in turn depends on the doping value  $N_D$  in the region below the ohmic contact. A two-region expression is used for the evaluation of  $\rho_c$  [15]:

$$\rho_c = \frac{10^{15}}{\mu_0 N_D} \quad \text{for } N_D \leq 4.7 \times 10^{17} \text{ atoms} \cdot \text{cm}^{-3} \quad (9a)$$

$$\rho_c = \frac{10^{15}}{\mu_0 4.7 \times 10^{17}} \quad \text{for } N_D > 4.7 \times 10^{17} \text{ atoms} \cdot \text{cm}^{-3}. \quad (9b)$$

The active layer resistance can be computed from the following expression.

$$R_{SL} = \int_0^{L_{SG}} \frac{dx}{Z q \int_0^{W(x)} N_D(y) \mu(y) dy}. \quad (10)$$

$W(x)$  is the active layer thickness and varies with  $x$  because of gate recess.

The evaluation of the intrinsic channel resistance  $R_i$  is based on a method described in the Appendix. The drain-source ( $C_{DS}$ ) and gate-drain ( $C_{GD}$ ) capacitances are finally calculated from the coupling between metal strips [16], [17]. The results of Fig. 2 show the excellent agreement between theoretical and experimental equivalent circuit data for an ion-implanted FET. The fringing capacitance ( $C_{GS_f}$ ) between source and gate metallizations is included in the calculation of  $C_{GS}$ . The experimental values were obtained by fitting calculated to measured  $S$ -parameter data in the frequency range of 2 to 18 GHz.

### C. Implantation Dose Influence on MESFET Performance

The simulated transistors had active layers realized by a two-step Si<sup>29</sup> ion implantation through a 500 Å thick Si<sub>3</sub>N<sub>4</sub> cap. The first implantation was done at an energy  $E_1 = 60$  keV and dose  $D_1 = 20 \times 10^{12} \text{ cm}^{-2}$ . These conditions were maintained the same throughout our investigations and resulted in a high doped N<sup>+</sup> surface layer. The second implantation had variable energies and doses. To examine the influence of dose, the energy was fixed at  $E_2 = 250$  keV and the dose was varied from  $D_2 = 3 \times 10^{12} \text{ cm}^{-2}$  to  $20 \times 10^{12} \text{ cm}^{-2}$ . For every  $D_2$  value, the recess depth ( $D_r$ ) is adjusted to yield a constant threshold voltage  $V_T = -3.6$  V (from  $D_r = 770$  Å for  $D_2 = 3.10^{12} \text{ cm}^{-2}$  to  $D_r = 2220$  Å for  $D_2 = 20 \times 10^{12} \text{ cm}^{-2}$ ).

The most relevant results are given in Fig. 3(a)–(c) and were obtained for constant  $V_{DS} = 3$  V,  $I_{DS} = 30$  mA bias. The intrinsic transconductance  $g_{m0}$  and gate capacitance  $C_{GS}$  are found to increase monotonically with dose (Fig. 3(a)). This is understood by noting that the recess depth was increased with the dose in order to maintain a constant threshold voltage. The thickness of the active layer below the gate was therefore thinner for high doses and resulted in higher  $g_{m0}$  and  $C_{GS}$ , as expected from the approximated expressions [18]

$$g_{m0} = \frac{\Delta I_{DS}}{\Delta V_{GS}} = \frac{\epsilon Z v_{sat}}{y_d} \quad (11a)$$

$$C_{GS} = \frac{\Delta Q}{\Delta V_{GS}} = \frac{\epsilon Z L}{y_d} \quad (11b)$$

where the depletion region thickness  $y_d$  is given by

$$y_d = \left[ \frac{2\epsilon}{qN_D} (V_{bi} - V_{GS}) \right]^{1/2}. \quad (11c)$$

Since the average donor concentration increases with dose,  $R_{DS}$  decreases with  $D_2$ , as shown in Fig. 3(b).  $R_i$  is also found to increase with dose. The parameters responsible for the latter change are, first, the higher doping  $N_D$  and, second, the smaller channel thickness  $a - y(x)$ . The nonuniform  $R_i(x)C_{GS}(x)$  transmission line under the gate has a resistance  $R_i(x)$  per unit length which is inversely proportional to the  $N_D(a - y(x))$  product. The rate at which  $N_D$  increases is obviously smaller than the corresponding thickness variations  $a - y(x)$  necessary to maintain a constant threshold. This results in an overall increase of  $R_i$  with dose.

The current cutoff frequency  $f_T = g_{m0}/2\pi C_{GS}$  remains almost constant with dose (Fig. 3(c)) since the  $g_{m0}$  and  $C_{GS}$  variations are very similar (see Fig. 3(a)). However, the largest  $f_{max}$  values (Fig. 3(c)) are obtained for small doses because  $R_i$  is, in this case, very small (see Fig. 3(b) and [18], [19]). The strong dependence of  $f_{max}$  on  $R_i$  demonstrates the need for very precise  $R_i$  evaluations to permit good understanding of the device microwave performance.

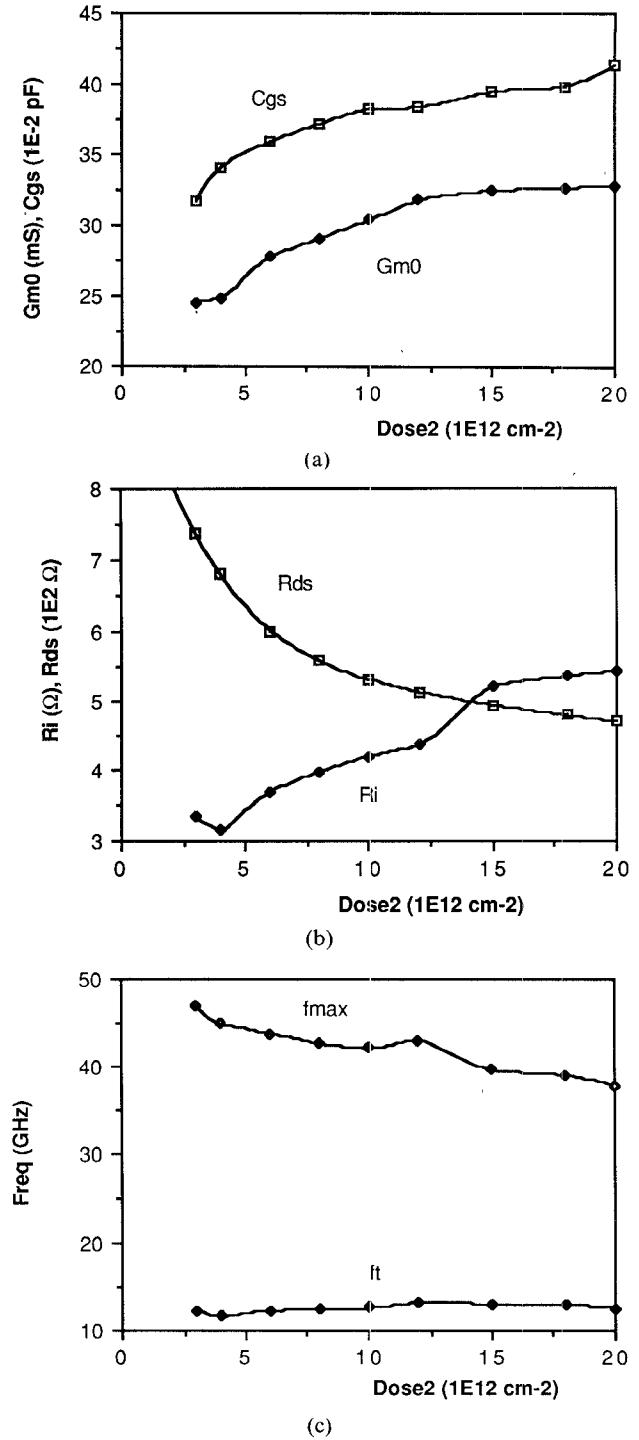


Fig. 3. Influence of implanted dose on: (a)  $C_{GS}$ ,  $g_{m0}$ , (b)  $R_{DS}$ ,  $R_i$ , and (c)  $f_T$  and  $f_{max}$  for an FET ( $300 \mu\text{m} \times 1 \mu\text{m}$ ,  $V_{DS} = 30$  mA) implanted with  $E_2 = 250$  keV.  $D_2$  varies between  $= 3 \times 10^{12} \text{ cm}^{-2}$  and  $20 \times 10^{12} \text{ cm}^{-2}$ , and gate recess was kept constant for  $V_T = -3.6$  V.

### D. Implantation Energy Influence on MESFET Performance

To investigate the influence of implantation energy, the dose of the second implants was kept constant at  $D_2 = 6 \times 10^{12} \text{ cm}^{-2}$ , and the energy varied from 100 keV to 400 keV. The recess depth was adjusted up to 2930 Å for  $E_2 = 400$  keV in order to yield a constant threshold voltage  $V_T = -3.6$  V. The bias conditions were  $V_{DS} = 3$  V,  $I_{DS} = 30$  mA.

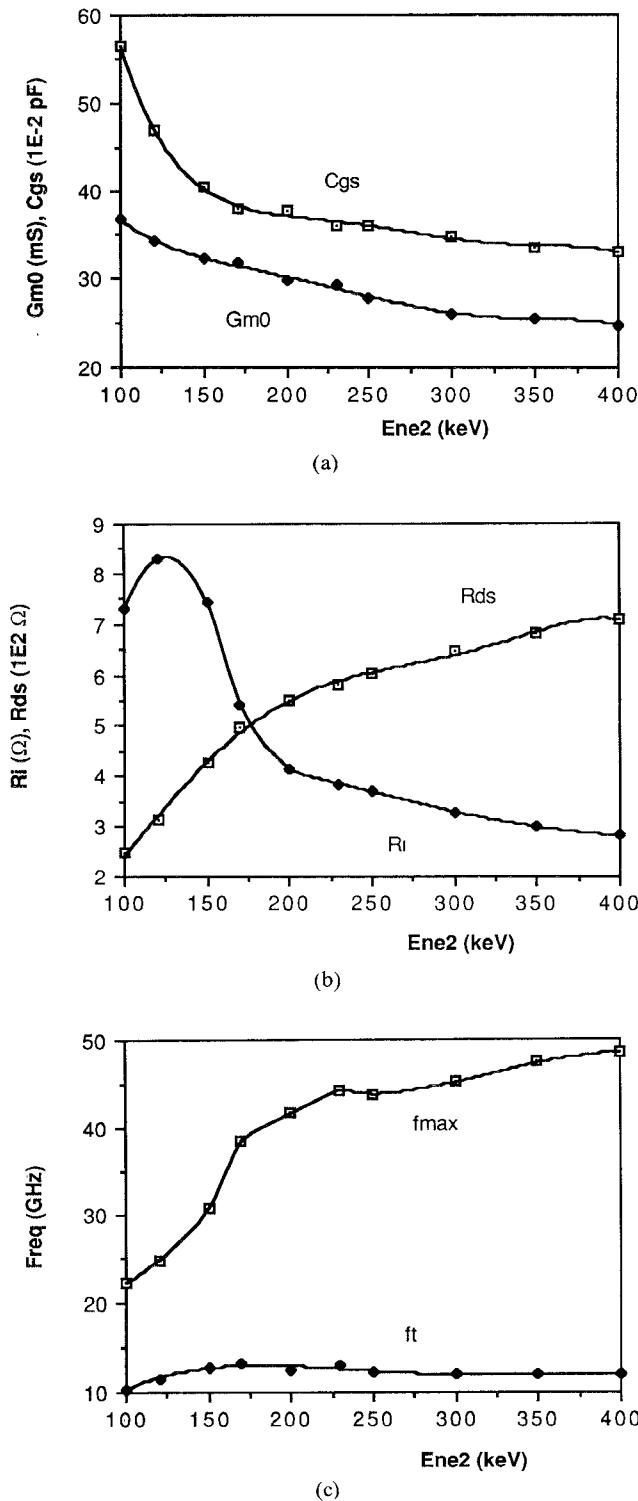


Fig. 4. Influence of implantation energy on (a)  $C_{GS}$ ,  $g_{m0}$ , (b)  $R_{DS}$ ,  $R_i$ , and (c)  $f_T$  and  $f_{max}$  for an FET ( $300 \mu\text{m} \times 1 \mu\text{m}$ ,  $V_{DS} = 30 \text{ mA}$ ) with constant implanted dose  $D_2 = 6 \times 10^{12} \text{ cm}^{-2}$  and energy  $E_2$  varying between 100 keV and 400 keV.  $V_T$  was kept constant, equal to  $-3.6 \text{ V}$ .

The active layer thickness increases with implantation energy. At higher energies and for recessed devices with the same  $V_T$ , the average donor concentration is lower in the channel below the gate. The associated depletion thickness is therefore larger. The  $g_{m0}$  and  $C_{GS}$  changes with implantation energy (Fig. 4(a)) can thus be explained by

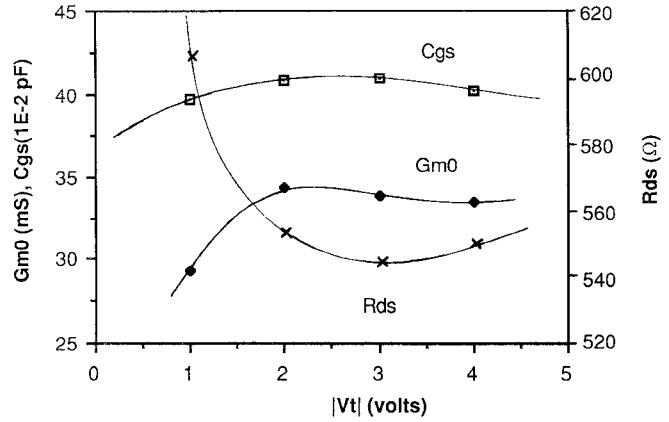


Fig. 5. Influence of threshold voltage ( $V_T$ ) on  $C_{GS}$ ,  $g_{m0}$ ,  $R_{DS}$  for an FET ( $300 \mu\text{m} \times 1 \mu\text{m}$ ,  $V_{DS} = 3 \text{ V}$ ,  $I_{DS} = I_{DSS}/2$ ), implanted with  $E_2 = 250 \text{ keV}$ ,  $D_2 = 6 \times 10^{12} \text{ cm}^{-2}$ . Different recess depths to obtain different  $V_T$  values.

depleted region extension. These results are in good agreement with the previously reported work of Trew [2].

The lower average donor concentrations are also responsible for the increase of  $R_{DS}$  with implantation energy (Fig. 4(b)).  $R_i$  decreases with  $E_2$  at values above 120 keV; at high energies, the low  $N_D$ 's have to be compensated by larger channel thicknesses to maintain  $V_T = \text{constant}$ , resulting therefore in smaller  $R_i$ 's.

It can again be noted that  $f_T$  is not very sensitive to energy variations of implantation (Fig. 4(c)). However,  $f_{max}$  is increased with energy due to both the reduction of the input resistance  $R_i$  and the increase of  $R_{DS}$ .

#### E. Gate Recess Depth Influence on MESFET Performance

For these investigations, the dose and energy of the second implants were maintained constant to  $D_2 = 6 \times 10^{12} \text{ cm}^{-2}$  and  $E_2 = 250 \text{ keV}$ . The bias conditions were  $V_{DS} = 3 \text{ V}$  and  $I_{DS} = I_{DSS}/2$ , where  $I_{DSS}$  is the drain-source saturation current at  $V_{GS} = 0 \text{ V}$ . (For instance,  $I_{DS} = 45 \text{ mA}$  when  $V_T = -3.6 \text{ V}$ .) The choice of these conditions allows device comparison since the absolute  $I_{DSS}$  value varies with recess. The recess depth was varied to obtain threshold voltages  $V_T$  between 1 and 5 V.

The variations of  $g_{m0}$ ,  $C_{GS}$ , and  $R_{DS}$  with recess are shown in Fig. 5. The high  $g_{m0}$  values at small recess depths (high  $V_T$ 's) are explained by the higher average carrier concentration when the complete implantation profile is available. When the recess depth is high, then the carrier distribution under the gate corresponds primarily to the tail of the profile and is therefore low, with the result of a small transconductance [18].

The slight increase of  $C_{GS}$  with  $V_T$  is again due to the higher carrier concentration  $n(y)$  at small recess depths. This is equivalent to smaller depletion thicknesses and consequently higher capacitance at large dopings.

$R_{DS}$  increases with recess depth (small  $V_T$ ) since the channel doping density is, in this case, relatively small. Similar behavior is also observed for  $R_i$ .

The frequencies  $f_T$  and  $f_{max}$  are increasing with  $V_T$  due to the faster changes for  $g_{m0}$  than  $C_{GS}$ .

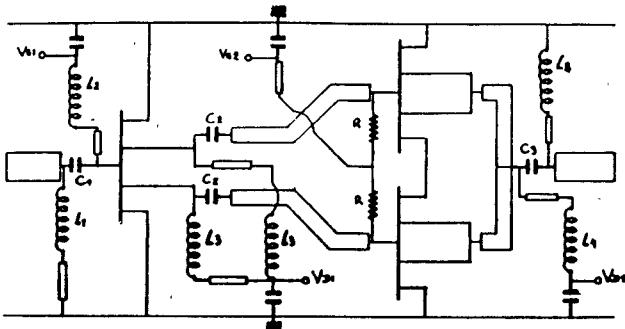


Fig. 6. Circuit diagram of the two-stage amplifier optimized by the SIMTEC FET model.

### III. AMPLIFIER DESIGN USING THE DEVELOPED MODEL

The FET model presented above has been successfully used in the design of microwave monolithic integrated GaAs amplifiers. It provides information regarding the influence that bias voltages, as well as implantation dose ( $D$ ), energy ( $E$ ), and threshold voltage ( $V_T$ ), have on the amplifier electrical characteristics. Small-signal gain ( $G$ ) and input and output  $SWR$  ( $SWR_{in}$  and  $SWR_{out}$ ) can thus be examined with respect to their sensitivity to material parameters ( $D, E$ ) and operating conditions. Circuit sensitivity to processing such as recess depth can also be considered by  $V_T$  simulations. The understanding of the above phenomena is of primary importance in the optimization of monolithic amplifiers since it makes it possible to design IC's which are tolerant to technology variations.

#### A. Amplifier Structure

Fig. 6 shows the circuit diagram of a two-stage monolithic amplifier designed with the tree approach [20], [21] and optimized with the model described in the previous sections. Each FET cell has a  $1\ \mu\text{m}$  gate length and a width of  $600\ \mu\text{m}$ , the total output periphery being consequently  $1200\ \mu\text{m}$ . Instead of using hybrid divider networks to feed the gates of a single transistor with large periphery, the tree approach uses the intrinsic isolation characteristics of FET's to divide the signals which have been amplified by the previous stages. It also uses several individual cells instead of a single large-periphery FET. The output combining network is made in lumped form and provides signal matching and dc biasing. These features allow the realization of compact power amplifier designs with no phasing or matching problems since individual cells have modest gate widths and do not suffer from impedance lowering. The amplifier of Fig. 6 was optimized for maximum small-signal gain and minimum input and output  $SWR$  over the frequency range of 6.0 to 7.0 GHz. It was intended for 10 percent bandwidth operation around the center frequency  $f = 6.5$  GHz. The gain roll-off over the operating bandwidth is 1.7 dB and reflects the best compromise for maximum gain and matched input-output. Large-signal characterization is outside the scope of this paper and can be studied by other models [22]; these take into consideration the change of the FET output resistance

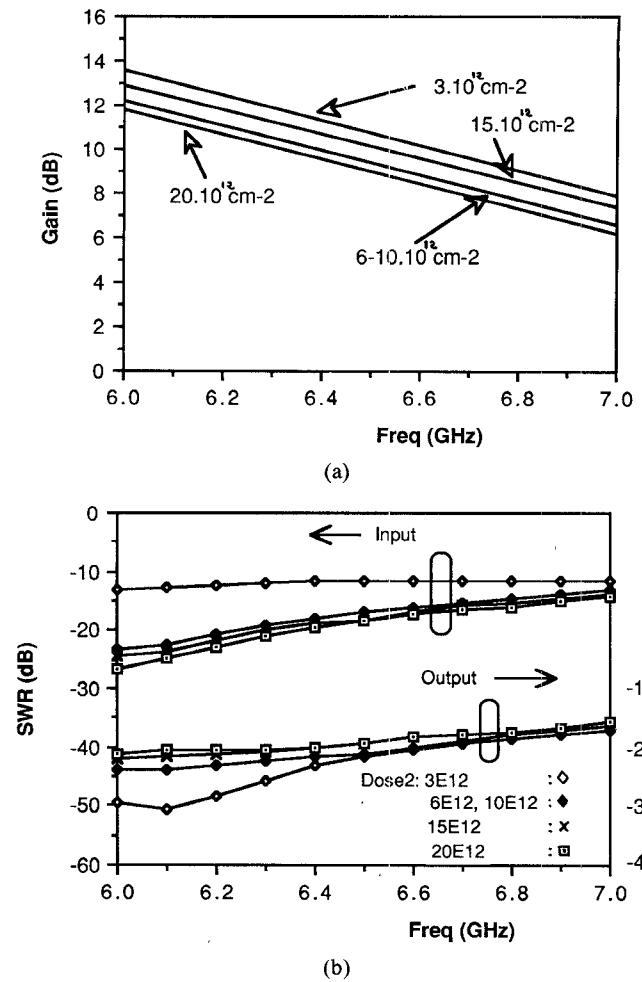


Fig. 7. Optimized amplifier performance and influence of implantation dose on (a) gain and (b) input and output  $SWR$ .

$R_{DS}$  with driving power and modify accordingly the previously optimized small-signal characteristics. The small-signal model used for the amplifier design is given in Fig. 2.

#### B. Dose Influence on Amplifier Performance

The optimized amplifier characteristics are shown in Fig. 7(a) and (b) together with the influence that dose variations have on the gain and  $SWR$ . In both cases, (a) and (b), the transistors have the same implantation energy ( $E_2 = 250$  keV) and threshold voltage ( $V_T = -3.6$  V), so that the influence of dose  $D_2$  can be studied alone. By decreasing the dose, the output resistance  $R_{DS}$  increases (Fig. 3(b)), and the input capacitance  $C_{GS}$  and transconductance  $g_m$  decrease (Fig. 3(a)) as outlined in the previous section. The changes in the values of the FET parameters with implanted dose can therefore be understood by considering the maximum unilateral gain of the FET, which is approximated by the following formula:

$$GMU \approx \frac{1}{16\pi^2} \frac{1}{f^2} \left( \frac{g_m}{C_{GS}} \right)^2 \frac{R_{DS}}{R_i + R_g + R_s} \quad (12)$$

where  $R_i$  and  $R_g$  are the intrinsic and metal resistance of the gate (see Fig. 2), and  $f$  is the frequency of operation. When the dose decreases from  $6 \times 10^{12}\ \text{cm}^{-2}$  to  $3 \times 10^{12}\ \text{cm}^{-2}$

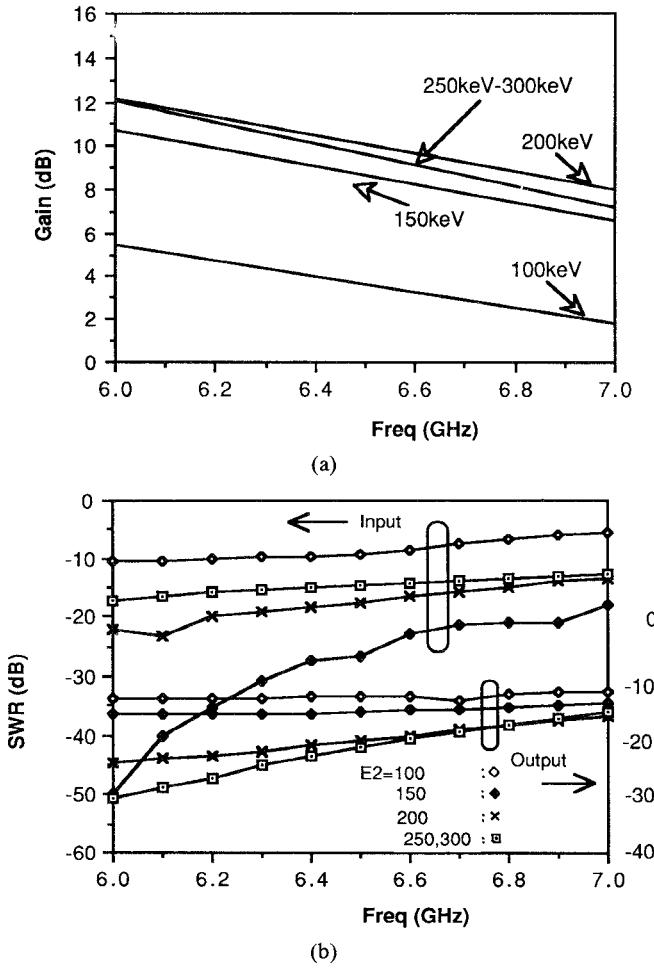


Fig. 8. Amplifier characteristics as a function of implantation energy.  
(a) Gain. (b) Input and output SWR.

$cm^{-2}$ ,  $R_{DS}$  changes faster than the  $g_m/C_{GS}$  ratio, resulting in FET's with improved GMU. This results in turn in higher amplifier gain with slightly improved output matching characteristics. The gate capacitance decrease is, however, so large that the input SWR of the amplifier is degraded. Although  $C_{GS}$  decreases rapidly for doses smaller than  $6 \times 10^{12} cm^{-2}$ , it shows relatively invariant characteristics for  $D_2$ 's up to  $20 \times 10^{12} cm^{-2}$ . This explains the large amplifier mismatch at  $3 \times 10^{12} cm^{-2}$  compared to  $20 \times 10^{12} cm^{-2}$ . The results suggest that the influence of implanted dose on circuit performance is very small.

### C. Energy Influence on Amplifier Performance

Implantation energy variations seem to have a more pronounced influence on circuit performance. The results of Fig. 8(a) and (b) show  $G$ ,  $SWR_{in}$ , and  $SWR_{out}$  curves of the two-stage amplifier for  $E_2$  energy variations between 100 and 300 keV (in all cases  $D_2 = 6 \times 10^{12} cm^{-2}$ ). Gate recess was again adjusted so that all FET's show the same  $V_T = -3.6$  V value. As the energy  $E_2$  decreases from 300 keV to 250 keV, the transconductance  $g_m$  and gate capacitance increase (Fig. 4(a)), but the output resistance  $R_{DS}$  decreases (Fig. 4(b)). These changes compensate each other and no substantial FET or amplifier performance variation

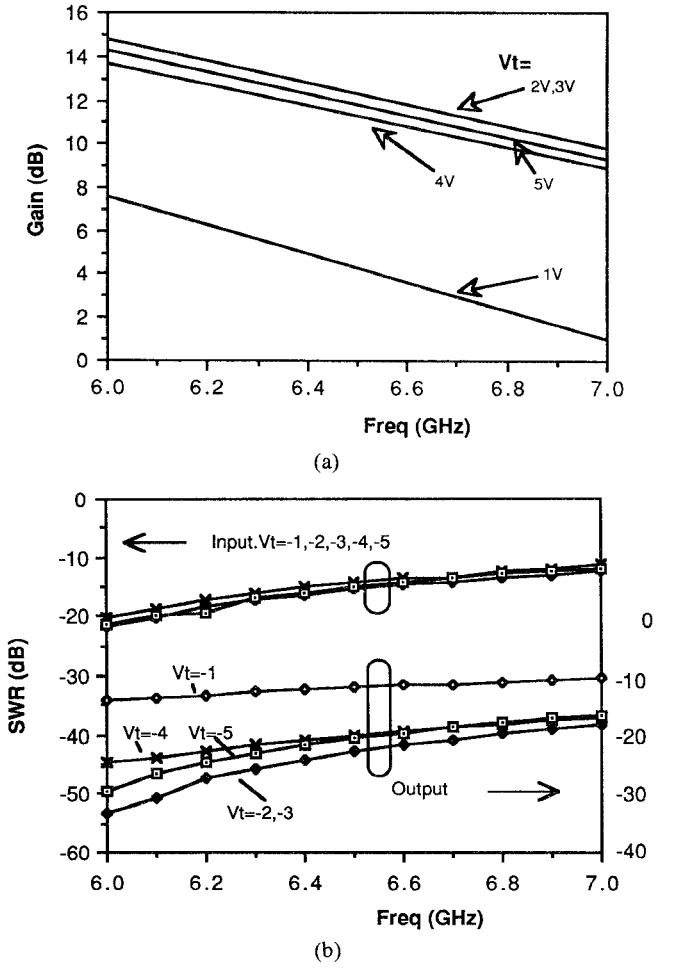


Fig. 9. Influence of threshold voltage  $V_T$  (recess) on the amplifier performance for a FET implanted with  $E_2 = 250$  keV,  $D_2 = 6 \times 10^{12} cm^{-2}$ . (a) Gain. (b) Input and output SWR.

is therefore observed between the nominal energy of 250 keV and  $E_2 = 300$  keV. For energies between 250 keV and 200 keV, the  $R_{DS}$  variation becomes slightly stronger than the  $g_m/C_{GS}$  change and the unilateral gain improvement of the FET produces a slight increase of amplifier gain at  $E_2 = 200$  keV. By further decreasing the energy, the output mismatch of the amplifier, caused by the rapid  $R_{DS}$  variations, results in steadily smaller gains and higher  $SWR_{out}$ . As far as the amplifier input is concerned, the slight increase of  $C_{GS}$  when decreasing the implanted energy from 300 keV to 150 keV seems to favor impedance matching. This is due to the input matching network being not perfectly optimized for the nominal energy  $E_2 = 250$  keV, since a compromise had to be made in optimizing simultaneously  $G$ ,  $SWR_{in}$ , and  $SWR_{out}$ . For even smaller energies, ( $E_2 = 100$  keV),  $C_{GS}$  is dramatically increased and causes pronounced amplifier mismatch.

### D. Recess Depth Influence on Amplifier Performance

The influence of recess on the amplifier performance is shown in Fig. 9(a) and (b). The nominal energy ( $E_2 = 250$  keV) and dose ( $D_2 = 6 \times 10^{12} cm^{-2}$ ) are used here for a FET with different threshold voltages  $V_T$ . The  $g_m$  and  $C_{GS}$  parameters increase initially with threshold voltages ( $V_T =$

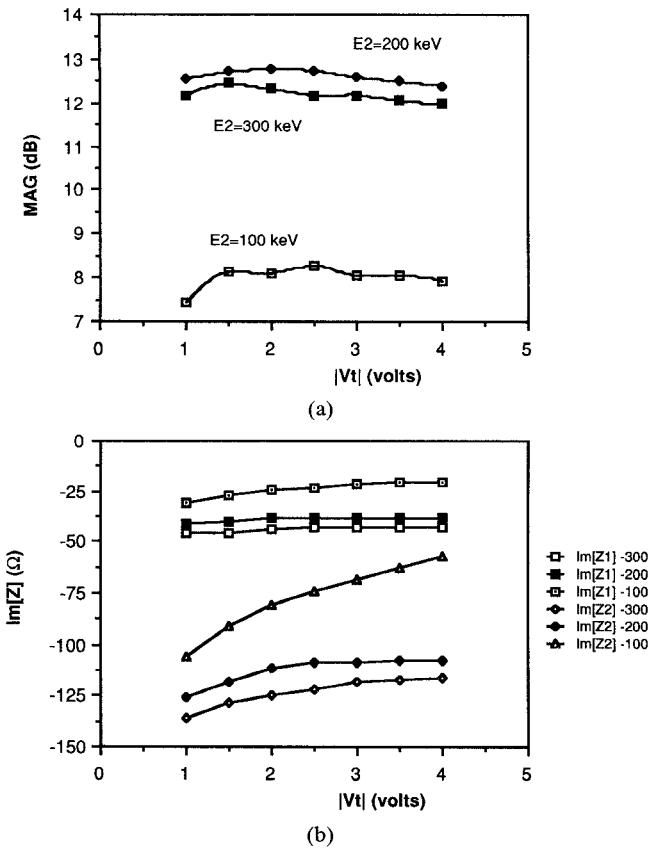


Fig. 10. (a) Maximum available gain (*MAG*) of MESFET ( $300 \mu\text{m} \times 1 \mu\text{m}$ ) at  $f = 6.5 \text{ GHz}$  as a function of recess (threshold voltage  $V_T$ ) and implantation energy ( $E_2$ ). Bias conditions are  $V_{DS} = 3 \text{ V}$ ,  $I_{DS} = I_{DSS}/2$ . (b) Imaginary part of input ( $Z_1$ ) and output ( $Z_2$ ) impedance of the same device.

1 V to 2 V) but remain almost constant above  $V_T = 2 \text{ V}$ . A similar but opposite tendency is observed for  $R_{DS}$ , which initially decreases for  $V_T = 1 \text{ V}$  to 2 V and then remains constant (Fig. 5). The rate of  $R_{DS}$  increase with recess is relatively faster than the  $C_{GS}$  decrease and this explains the greater dependence of output rather than input matching on  $V_T$  variations. The large output mismatch created at low  $V_T$ 's is partly responsible for the smaller amplifier gain.

#### E. Influence of Profile on Maximum Available Gain (*MAG*)

Sections III-B to III-D presented the influence of dose, energy, and recess on the amplifier characteristics. The validity of these results was confirmed by investigating the amplifier and FET *MAG* as a function of  $V_T$  and  $E_2$ . The data obtained in this way depend entirely on profile characteristics since perfect matching is guaranteed at both the input and output. Fig. 10 shows such results for a MESFET at 6.5 GHz. Maximum gain (Fig. 10(a)) is obtained again for  $E_2 = 200 \text{ keV}$  (see Section III-C) while no significant *MAG* dependence on  $V_T$  is observed. Similar results were obtained for the amplifier by perfect matching of input and output and interstage optimization.

The profile characteristics should not, however, be considered only in terms of their influence on gain. As explained in Section II, III-C and III-D, the  $E_2, V_T$  choice

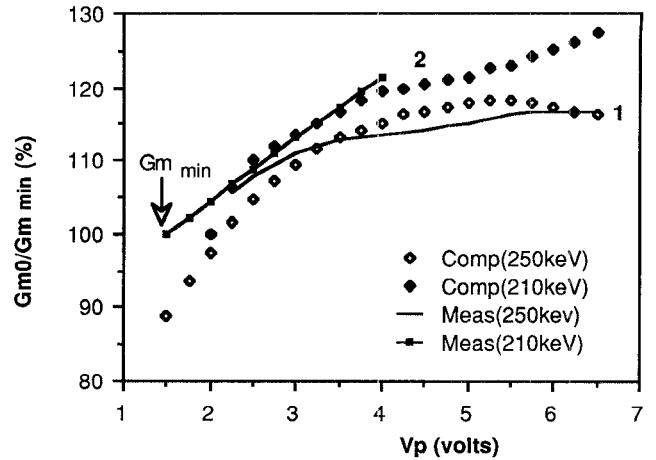


Fig. 11. Computed (scattered points) and measured (solid lines) normalized transconductance ( $G_{m0}/G_{m\min}$ ) versus different pinch-off voltages as obtained by varying the gate recess depth ( $G_{m\min} = 90 \text{ mS/mm}$  is the smallest transconductance measured). Two wafers realized by double implantation of Si are investigated. Wafer 1, 60 keV–250 keV. Wafer 2, 60 keV–210 keV.

also has an effect on the input/output device impedance. For amplifier design this signifies reoptimization of matching in order to benefit from the new gain characteristics. Fig. 10(b) shows the influence of  $V_T, E_2$  on the imaginary part of MESFET input and output impedance ( $Z_1, Z_2$ ).  $Z_1$  variations are primarily attributed to  $C_{GS}$ , while  $Z_2$  changes are due to  $R_{DS}$  and  $g_{m0}$  [18]. The resistive part of  $Z_1, Z_2$  remains almost constant with  $V_T, E_2$ . Process tolerance will also be subjected to the simultaneous effect of gain and matching.

## IV. EXPERIMENTAL RESULTS

### A. Recess Depth Influence on Device Transconductance

Two series of MESFET's with ( $1 \mu\text{m} \times 150 \mu\text{m}$ ) gates have been realized by double Si implantation through 500 Å of  $\text{Si}_3\text{N}_4$  (series 1:  $E_1 = 60 \text{ keV}$  followed by  $E_2 = 250 \text{ keV}$  and series 2:  $E_1 = 60 \text{ keV}$  followed by  $E_2 = 210 \text{ keV}$ ). For both series, the dose corresponding to the 60 keV energy was  $D_1 = 20 \times 10^{12} \text{ cm}^{-2}$ . The dose of the second implantation step was  $D_2 = 6 \times 10^{12} \text{ cm}^{-2}$ . Different recess depths were realized on the same wafer in order to obtain devices with different pinch-off voltages ranging from 2 V to 6 V. For each device the transconductance  $g_{m0}$  was measured. The comparison with the theoretical values is shown in Fig. 11, where the relative transconductance is plotted versus the pinch-off voltage.

When the implantation energy increases (series 1)  $g_{m0}$  decreases and its dependence on recess depth becomes smaller. The simulations are in good agreement with these experimental results and indicate that the simulator SIMTEC can account equally well for recess depth and doping profile shape.

It is also interesting to note that for given implantation energy,  $g_{m0}$  is found to increase with  $V_T$  at a much faster rate for low, rather than high, implantation energies. This suggests the use of low energies for applications where the

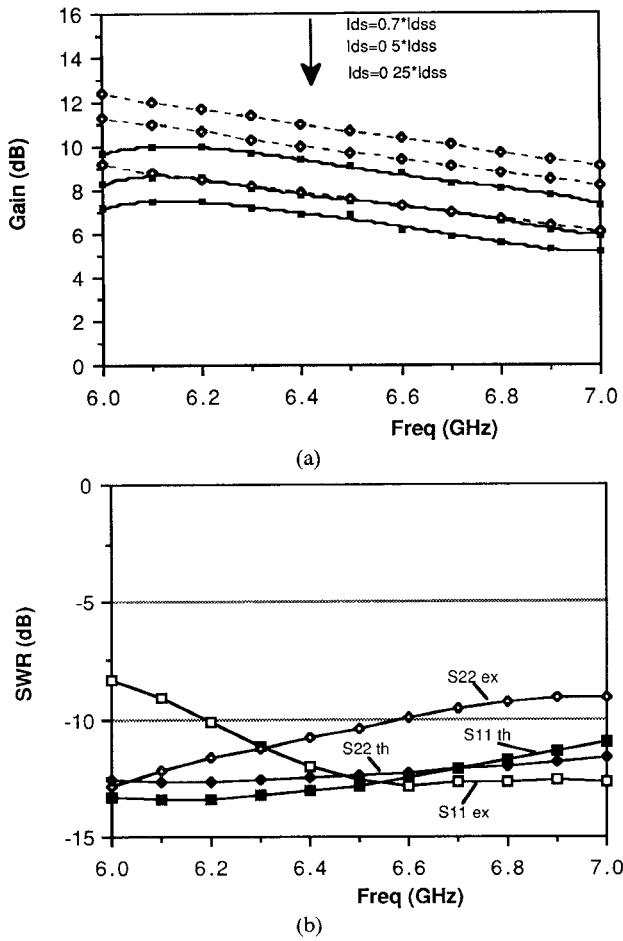


Fig. 12. Experimental (solid lines) and theoretical (dashed lines) microwave performance of the amplifier with  $E_2 = 250$  keV,  $D_2 = 6 \times 10^{12}$  cm $^{-2}$  active layers and  $V_T = -3.6$  V. The gain (a) is plotted as a function of the drain-source current  $I_{DS} = 0.7I_{DSS}$ ,  $0.5I_{DSS}$  and  $0.25I_{DSS}$  ( $I_{DSS}$  = saturation current). The input and output SWR (b) are plotted for  $I_{DS} = 0.5I_{DSS}$ .

$g_{m0}$  improvement is absolutely essential. It also indicates that recessing and, therefore, process variations will have less influence on the transistor  $g_m$  if high energy profiles are chosen [7], [9].

#### B. Bias Influence on Amplifier Performance

The influence of bias on the performance of the two-stage tree amplifier is finally shown in Fig. 12(a) and (b). This study allowed also a first comparison between experimental and theoretical results for the amplifier. The simulations were made using the EESOF Package Touchstone [23] and include consideration of: (i) all "parasitic" connections between MMIC components, bias pads, and ground, (ii) the external biasing capacitances (50 pF) necessary for better RF grounds, and (iii) transitions to the tapered coplanar lines of the test package. The influence of other effects, such as parasitic capacitances of lumped elements and losses of transmission lines and components, have also been considered.<sup>1</sup>

<sup>1</sup>The final circuit file, together with other relevant information, is not given due to limited space, but it will be supplied by the authors upon request.

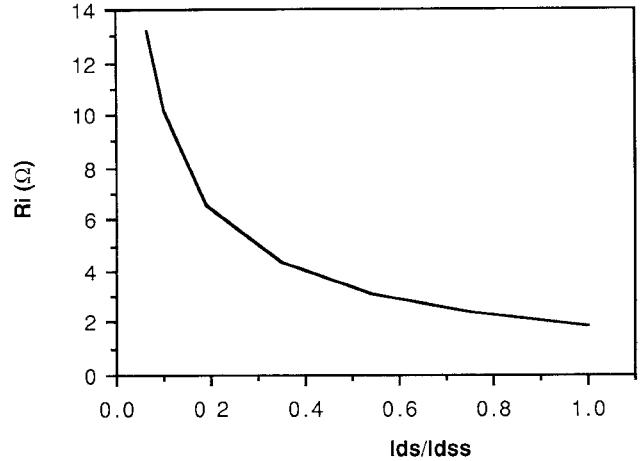


Fig. 13. Computer variations of  $R_i$  versus  $I_{DS}$ . The FET is the same as in Fig. 2.

In Fig. 12(a), the amplifier gain  $G$  is plotted over the frequency band of 6 to 7 GHz for bias operating conditions corresponding to  $0.7I_{DSS}$ ,  $0.5I_{DSS}$ , and  $0.25I_{DSS}$ . The agreement between theory and experiment is very satisfactory (between 1 and 2.5 dB) in spite of the fact that (i) the experimental results are average values among five chips of the wafer and the relative dispersion of the measurements is  $\Delta G = \pm 2.0$  dB; (ii) the  $g_m$  values used in the circuit simulations (170 mS/mm at  $I_{DSS}$ ) were calculated by the presented MESFET model; these are in good agreement with measured values of carefully selected discrete devices (see agreement in Fig. 2), but are, in fact, slightly higher than the measured integrated transistors; and (iii) the measurements are made at the plane of the package connectors since no monolithic calibration standards were available. The precision measurement technique reported earlier [24] could not be applied because the amplifier was designed for use in a package. As shown by both the theoretical and the experimental results, a change in the drain-source current from  $0.25I_{DSS}$  to  $0.7I_{DSS}$  produces an overall gain improvement. This can be understood by considering the  $g_m$ ,  $C_{GS}$  increase, as well as the  $R_{DS}$ ,  $R_i$  variations: (i)  $g_m$  varies with  $I_{DS}$  faster than  $C_{GS}$ ; (ii)  $R_{DS}$  decreases with  $I_{DS}$  but changes at about the same rate as  $C_{GS}$ ; (iii) the intrinsic gate resistance  $R_i$  reduces by as much as 54 percent when biasing at  $0.7I_{DSS}$  instead of  $0.25I_{DSS}$  (Fig. 13). In fact, the rate of  $R_i$  change with  $I_{DS}$  is much faster than for  $C_{GS}$  and, combined with the variations of  $g_m$  and  $R_{DS}$ , results in a larger FET and amplifier gain.

A comparison between theoretical and experimental results concerning the input and output SWR is given in Fig. 12(b), for a biasing  $V_{GS} = 3$  V,  $I_{DS} = 0.5I_{DSS}$ . The agreement is excellent and the differences of input SWR in the range of 6–6.5 GHz can be accounted for by the lack of exact calibration, as discussed earlier.

#### V. CONCLUSIONS

A device simulator called SIMTEC has been presented and applied to ion-implanted MESFET's. Although based on a certain number of assumptions, like most models of

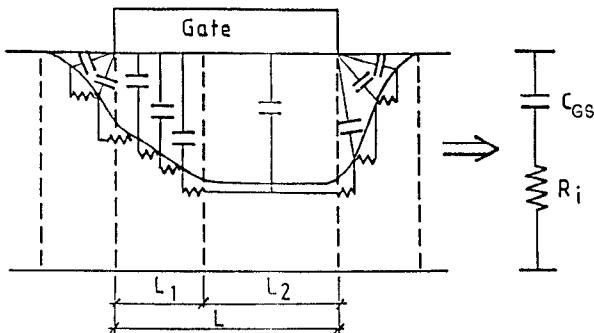


Fig. 14. Transmission line model used for the calculation of the input network  $R_i - C_{GS}$ . Only eight cells are presented, each corresponding to a cascaded uniform line, but 20 are used in the simulation.

this kind, it proved to be efficient for the technology and process simulation of individual devices and multistage MMIC amplifiers. A good qualitative and quantitative understanding of MESFET and MMIC amplifier influence by implantation dose, energy, recess, and bias conditions was possible. The amplifier gain agreement between theory and experiment is very satisfactory and *SWR* agreement is also very good. SIMTEC is therefore suitable for computer-assisted optimization of MMIC's with respect to processing and sensitivity to technology dispersions.

The results obtained for a two-stage MMIC amplifier show that the gain and input *SWR* are not very sensitive to the implanted dose as long as this exceeds  $6 \times 10^{12} \text{ cm}^{-2}$ . Implantation energies larger than 200 keV led to smaller gain and *SWR* dispersion. Recess depths should finally be small ( $V_T > 2 \text{ V}$ ) for the microwave performance to become insensitive to recess variations.

An improvement of MMIC performance and processing yield can therefore be expected by the use of the developed simulator.

## APPENDIX

The input impedance of a common-source MESFET corresponds to that of a nonuniform  $R - C$  transmission line (Fig. 14) made of (i) the distributed capacitance of the space charge below the gate (or below the surface of the uncovered active layer) and (ii) of the distributed resistance of the neutral channel. The transmission line is loaded at its drain end by the saturated section of the channel and can therefore be considered as open-circuited.

The computations of the nonuniform line input impedance were performed by cascading about 20 different  $R - C$  transmission lines, each of them being considered uniform (Fig. 14). The input impedance of all these cascaded lines can therefore be easily obtained and the real part of it is equal to the intrinsic channel resistance  $R_i$ .

$C_{GS}$  can be obtained either from the imaginary part of the impedance or from the incremental charge of the stored charge below the gate corresponding to a variation of intrinsic gate voltage. Both techniques yield similar results.

Finally, the edge effects are implicitly included in the present calculations since the space charge below the

surface (due to Fermi level pinning) and at the gate edges is considered.

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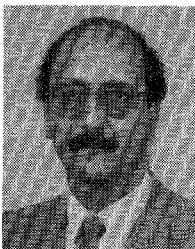
The simulations reported in this work were performed at the computer facilities at the University of Michigan, and the continuous help and support of the computing staff, in particular C. Nicholas, are greatly acknowledged.

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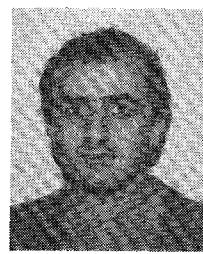
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Dr. Graffeuil has authored or coauthored over 30 technical papers and three books.